



Cite this: *Nanoscale*, 2018, **10**, 2363

## Reduced sintering of mass-selected Au clusters on SiO<sub>2</sub> by alloying with Ti: an aberration-corrected STEM and computational study†

Yubiao Niu,<sup>a</sup> Philomena Schlexer,<sup>b</sup> Bela Sebok,<sup>c</sup> Ib Chorkendorff,<sup>id c</sup>  
Gianfranco Pacchioni<sup>d</sup> and Richard E. Palmer<sup>id \*d</sup>

Au nanoparticles represent the most remarkable example of a size effect in heterogeneous catalysis. However, a major issue hindering the use of Au nanoparticles in technological applications is their rapid sintering. We explore the potential of stabilizing Au nanoclusters on SiO<sub>2</sub> by alloying them with a reactive metal, Ti. Mass-selected Au/Ti clusters (400 000 amu) and Au<sub>2057</sub> clusters (405 229 amu) were produced with a magnetron sputtering, gas condensation cluster beam source in conjunction with a lateral time-of-flight mass filter, deposited onto a silica support and characterised by XPS and LEIS. The sintering dynamics of mass-selected Au and Au/Ti alloy nanoclusters were investigated in real space and real time with atomic resolution aberration-corrected HAADF-STEM imaging, supported by model DFT calculations. A strong anchoring effect was revealed in the case of the Au/Ti clusters, because of a much increased local interaction with the support (by a factor 5 in the simulations), which strongly inhibits sintering, especially when the clusters are more than ~0.60 nm apart. Heating the clusters at 100 °C for 1 h in a mixture of O<sub>2</sub> and CO, to simulate CO oxidation conditions, led to some segregation in the Au/Ti clusters, but in line with the model computational investigation, Au atoms were still present on the surface. Thus size-selected, deposited nanoalloy Au/Ti clusters appear to be promising candidates for sustainable gold-based nanocatalysis.

Received 24th August 2017,  
Accepted 6th October 2017

DOI: 10.1039/c7nr06323g

rsc.li/nanoscale

### 1. Introduction

Gold (Au) is considered to be the noblest of all metals, long regarded as chemically rather inert<sup>1</sup> and thus (more or less) catalytically inactive. Nevertheless a report from as long ago as 1906 states that a hot Au mesh can facilitate the combination of O<sub>2</sub> and H<sub>2</sub>,<sup>2</sup> while a paper from 1925 indicates that Au is capable of catalysing the oxidation of CO to CO<sub>2</sub>.<sup>3</sup> But since Au is fairly expensive and the reported reaction rates were lower compared with other metals, Au was not considered as a viable catalyst until Sennewald *et al.*<sup>4</sup> found that a Pd–Au catalyst was active for the oxidative acetoxylation of ethylene to vinyl acetate in 1965 and then Bond and Sermon<sup>5</sup> demonstrated the activity

of supported Au particles with diameters of 2–100 nm for the hydrogenation of 1-pentene at 100 °C in 1973. Later, in the 1980s, Haruta *et al.*<sup>6,7</sup> discovered that Au nanoparticles, with diameters of 3–5 nm hold up on reducible supports, were exceptionally active for CO oxidation even at –70 °C, and Hutchings<sup>8</sup> demonstrated the catalytic activity of Au<sup>3+</sup> in the hydrochlorination of acetylene to vinyl chloride. Since then, Au has attracted growing attention as a catalyst and become the preminent example of a size effect in heterogeneous catalysis.<sup>9–11</sup>

As CO oxidation is one of the simplest reactions which Au can catalyse, and CO is frequently used as a probe molecule in surface science, CO oxidation on TiO<sub>2</sub> supported Au nanoparticles has become a prototypical model system in Au catalysis,<sup>12–15</sup> and numerous studies have been reported regarding the reaction mechanisms,<sup>16–18</sup> active sites,<sup>11,19–21</sup> and active species.<sup>22,23</sup> Looking at the experimentally measured catalytic activities of Au nanoparticles on different supports in the CO oxidation reaction, two main conclusions can be drawn. The Au nanoparticles are not active above a size of approx. 5 nm and on some oxide supports, such as SiO<sub>2</sub>, they are not active at all, regardless of size.<sup>11</sup>

Although supported Au nanoparticles have genuine potential in technological applications, a major issue hindering

<sup>a</sup>Nanoscale Physics Research Laboratory, School of Physics and Astronomy, University of Birmingham, B15 2TT Birmingham, UK

<sup>b</sup>Dipartimento di Scienza dei Materiali, Università di Milano-Bicocca, 20125 Milano, Italy

<sup>c</sup>Department of Physics, SurfCat, Technical University of Denmark, DK-2800 Kgs. Lyngby, Denmark

<sup>d</sup>College of Engineering, Swansea University, Bay Campus, Swansea, SA1 8EN, UK.  
E-mail: r.e.palmer@swansea.ac.uk

†Electronic supplementary information (ESI) available. See DOI: 10.1039/c7nr06323g



their implementation is rapid sintering. It has been shown that Au atoms are mobile and can migrate to form 3D islands on a TiO<sub>2</sub> surface even at 150–160 K.<sup>24,25</sup> Also on silica, Au particles tend to sinter quickly, unless defects are present on the silica surface to which the Au particles bind more strongly.<sup>26</sup> However, the defects have to be present on the surface before the cluster deposition and are unlikely to be formed *in situ*.<sup>27</sup> According to the work of Yang *et al.*,<sup>28</sup> Au clusters on TiO<sub>2</sub> sinter by Ostwald ripening between 300 K and 410 K and the sintering is accelerated by the presence of a mixture of CO and O<sub>2</sub>. The sintering modes of mass-selected Au nanoclusters deposited on amorphous carbon were studied by Hu *et al.* They found that the sintering process is size-dependent; Au<sub>561±13</sub> and Au<sub>923±20</sub> clusters exhibit Ostwald ripening, whereas Au<sub>2057±45</sub> ripens through cluster diffusion and coalescence (Smoluchowski ripening).<sup>29</sup>

The stabilisation of supported nanoparticles against sintering has attracted significant research effort.<sup>30,31</sup> Two pathways are usually described: alloying the particles, or encapsulating them within an oxide or organic shell. The latter may hinder the use of nanoparticles in catalytic applications because of surface blocking, but alloying may provide a route to nanoparticles which are both stable and active. In particular, the investigations of mass-selected alloy nanoparticles synthesised in the gas phase have already shown the unique properties of such materials. Examples include increased activity in the oxygen reduction reaction (ORR) for case of Pt<sub>x</sub>Y and Pt<sub>x</sub>Gd nanoparticles as a result of strain<sup>32,33</sup> and the dynamic behavior of CuZn nanoparticles under methanol synthesis reaction conditions.<sup>34</sup> With regards to sintering, it has been shown that alloying Au with Ir, Cu or Ag increases the stability of chemically synthesized nanoparticles on oxide supports.<sup>35–41</sup> However, studies of the effects of alloying on the stability of mass-selected nanoparticles against sintering have been lacking to date.

Here we show that the sintering rate of mass-selected Au/Ti alloy nanoclusters, synthesised in a magnetron sputtering, gas condensation cluster beam source<sup>42,43</sup> is much lower than that of mass-selected pure Au nanoclusters on silica. The composition of the clusters was characterised by X-ray Photoelectron Spectroscopy (XPS) and Low Energy Ion Scattering (LEIS) while aberration-corrected Scanning Transmission Electron Microscopy (STEM) was used to investigate the sintering process *via* direct real space imaging with atomic resolution. Complementary *ab initio* calculations confirm stronger binding between alloyed Au/Ti clusters and the SiO<sub>2</sub> surface compared with pure Au clusters.

## 2. Experimental and computational details

### 2.1 Experimental details

Au<sub>2057</sub> (405 229 amu) and Au/Ti (400 000 amu) nanoclusters were produced with a DC magnetron sputtering, gas condensation cluster beam source<sup>43</sup> from Au (PI-KEM, 99.99%) and

Au/Ti (PI-KEM, 50/50 at%, 99.99%) targets, respectively. The clusters with a specific mass were selected by a lateral time-of-flight mass filter<sup>42</sup> before deposition onto the support with a kinetic energy of approx. 0.5 eV per atom. The mass-selected clusters were deposited onto thin silica films suspended on copper TEM grids (EMS, USA) and onto solid Si substrates terminated with a 50 nm thermally grown SiO<sub>2</sub> layer. These substrates were equipped with heater elements and a Pt resistance temperature detector (RTD) evaporated on their back side. The loading of all the samples was approx. 10% (projected surface area coverage), *i.e.*, approx. 10% of the surface covered by clusters. Before deposition the Si/SiO<sub>2</sub> substrates were sputtered by an Ar<sup>+</sup> beam from the cluster source, with a beam current of 1 nA for 500 s (at 1000 eV Ar<sup>+</sup>) in order to clean the surface and create defects to act as anchor points for the clusters. After deposition on the TEM supports, the clusters were imaged with a 200 kV aberration-corrected JEOL 2100F Scanning Transmission Electron Microscopy (STEM) in high-angle annular dark field (HAADF) mode after transfer under atmospheric conditions. HAADF images were acquired with inner and outer detector angles of 62 and 164 mrad (camera length 10 cm), and probe convergence angle of 19 mrad. The Si/SiO<sub>2</sub> samples were characterized with XPS (non-monochromatized Mg K $\alpha$  line from a SPECS XR50 X-ray gun, using an OMICRON NanoSAM 7 channel energy analyzer) and LEIS (1 KeV He<sup>+</sup> from an Omicron ISE100 ion gun using the same energy analyzer), again after transfer under atmospheric conditions. The silicon supports were heated in 1 bar of O<sub>2</sub>:CO = 4:1 to approx. 370 K (100 °C) in order to simulate the effect of CO oxidation conditions. We estimate the error in temperature here to be around  $\pm 15$  K. After heating the supports were characterized again with XPS and LEIS to investigate the surface composition of the clusters and the effect of the heating.

### 2.2 Computational details

We performed spin polarized density functional theory (DFT) calculations with the Vienna Ab Initio Simulation Package (VASP 5.2)<sup>44–47</sup> using the Perdew, Burke and Ernzerhof (PBE) exchange–correlation functional.<sup>48,49</sup> Electron-ion interactions were described *via* the projector augmented wave (PAW) method.<sup>50,51</sup> O(2s, 2p), Si(3s, 3p), Au(5d, 6s) and Ti(3d, 4s, 4p) states were treated explicitly. For electronic relaxations, we used the blocked Davidson iteration scheme.<sup>52,53</sup> Full geometric structure optimizations were performed, allowing all ions to relax until forces were smaller than  $|0.01|$  eV Å<sup>-1</sup>. The lattice parameters were kept fixed during structure optimizations.

To represent the silica surface, we used a periodic fully hydroxylated  $\alpha$ -quartz (001) surface, as described in our previous study.<sup>54</sup> A (3  $\times$  3) surface super cell was used with lattice parameters of  $a = b = 15.11$  Å and  $\gamma = 120^\circ$  and with a slab thickness of 9 layers of [SiO<sub>4</sub>] tetrahedra. The slabs were separated by more than 20 Å of vacuum to ensure enough space for the metal clusters. Wave functions were expanded in a plane wave basis up to a kinetic energy of 400 eV. A  $\Gamma$ -centred



K-point grid in the Monkhorst–Pack scheme<sup>55</sup> was used, which was set to the  $\Gamma$ -point.

As dispersion forces can be important for the cluster-support interaction, we applied the semi-empirical dispersion correction proposed by Grimme,<sup>56</sup> known as the DFT-D2 approach. As the DFT-D2 approach is assumed to overestimate the dispersion in oxides, we changed the  $C_6$  and  $R_0$  parameters as suggested by Tosoni and Sauer.<sup>57</sup> The resulting approach is called DFT-D2'. Atomic charges  $q$  were determined *via* the Bader decomposition scheme.<sup>58–60</sup> The adsorption energies of the clusters,  $E_{\text{ADS}}$ , are defined in eqn (1), where  $E(\text{Au}_x\text{Ti}_y)$  are the metal clusters in the gas-phase and  $S$  is the support.

$$E_{\text{ADS}}(\text{Au}_x\text{Ti}_y/\text{S}) = E(\text{Au}_x\text{Ti}_y/\text{S}) - E(\text{Au}_x\text{Ti}_y) - E(\text{S}) \quad (1)$$

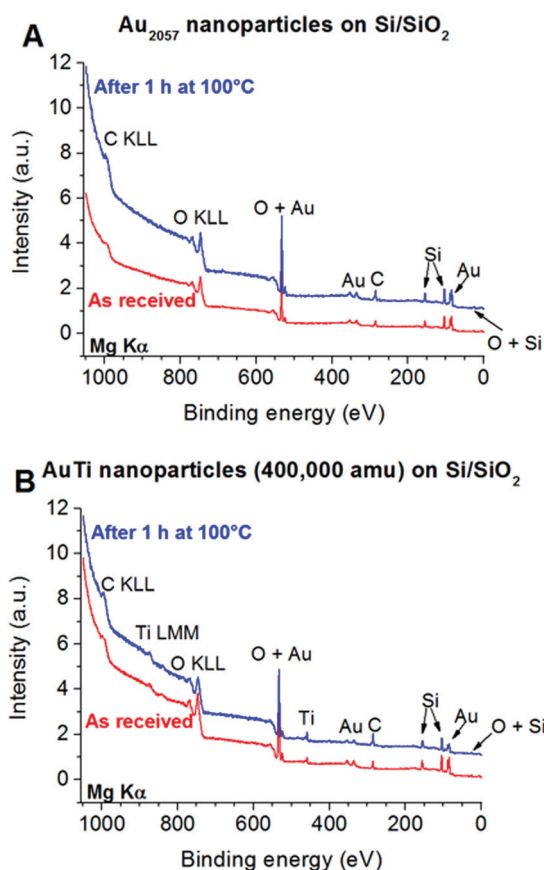
### 3. Results and discussion

#### 3.1 XPS and LEIS of the nanoclusters

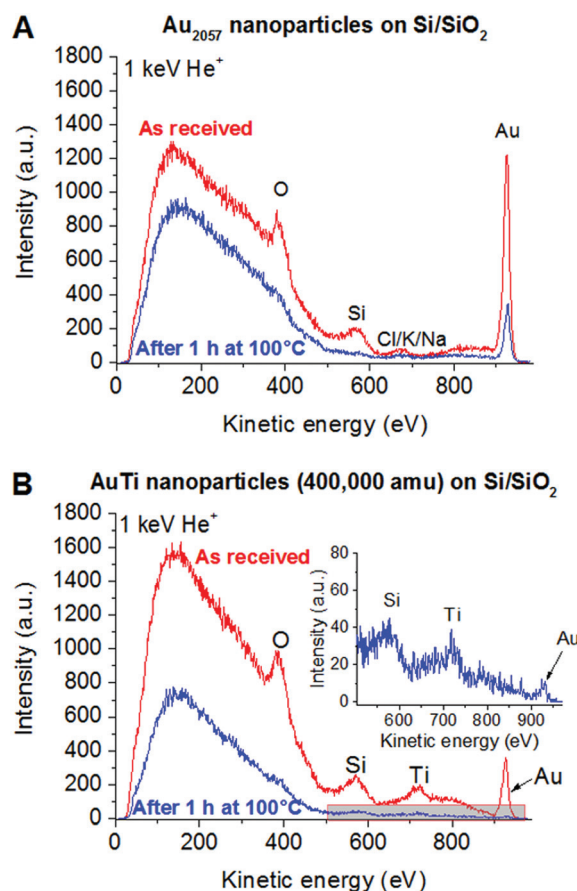
XPS was used to provide information on the composition of the deposited clusters. Typical XPS spectra are shown in Fig. 1 for both the  $\text{Au}_{2057}$  and  $\text{Au}/\text{Ti}$  nanoclusters deposited on

$\text{Si}/\text{SiO}_2$  supports, as received after transport under atmospheric conditions, and also after heating in a mixture of  $\text{O}_2$  and  $\text{CO}$  in order to simulate the effect of  $\text{CO}$  oxidation conditions. In all samples the Si in  $\text{SiO}_2$  binding state, originating from the substrate, can be detected. Traces of carbonaceous contaminants from atmospheric exposure are also visible in the spectra. In case of the  $\text{Au}_{2057}$  nanocluster sample, no significant changes are visible in the XPS spectra taken before and after heating in  $\text{O}_2/\text{CO}$  (Fig. 1A). In the case of the  $\text{Au}/\text{Ti}$  nanoclusters, the relative intensity of the Au signal compared with the Si signal (normalized peak size) in the spectrum decreases after heating. This is consistent with a segregation process in which more Ti is drawn to the cluster surface, which is underlined by the change of the Au : Ti ratio determined from the spectra. The as received clusters show a surface Au : Ti ratio of 60 : 40, while after heating the ratio is approx. 40 : 60. Some Ti is present on the surface in the form of  $\text{TiO}_2$  regardless of the heating (see detailed spectra in ESI, Fig. S1†).

In order to further investigate the surface of the samples, LEIS spectra were recorded (Fig. 2) characterizing the outermost layer of the sample surface. In both Au and  $\text{Au}/\text{Ti}$  cluster samples, O, Si and Au are detectable on the surface before heating, as expected for Au-containing clusters deposited on a



**Fig. 1** XPS spectra of  $\text{Au}_{2057}$  (A) and  $\text{Au}/\text{Ti}$  (B) nanoclusters deposited on  $\text{Si}/\text{SiO}_2$  slabs before and after  $100^\circ\text{C}$  for 1 h in 1 bar of  $\text{O}_2:\text{CO} = 4:1$  mixture. Spectra were calibrated for Au  $4f_{7/2}$  line to appear at binding energy of 84 eV and were normalized with the area of the Si 2s peak. The plots were offset in y direction for better visibility.



**Fig. 2** LEIS spectra of  $\text{Au}_{2057}$  (A) and  $\text{Au}/\text{Ti}$  (B) nanoclusters deposited on  $\text{Si}/\text{SiO}_2$  slabs before and after (inset)  $100^\circ\text{C}$  for 1 h in 1 bar of  $\text{O}_2:\text{CO} = 4:1$  mixture taken with 1 keV  $\text{He}^+$  ions.



SiO<sub>2</sub> surface. In case of the sample with Au<sub>2057</sub> clusters (Fig. 2A), small additional contaminant(s) can be detected having a mass of approx. 39 amu (Cl, K or Na), which could originate from the handling of the sample. No such contaminant could be detected in case of the sample with Au/Ti clusters, since Ti is clearly visible in the same region (Fig. 2B). After heating the samples, the Si and O peaks disappear from the spectra. As XPS confirms the presence of SiO<sub>2</sub>, this disappearance is most probably the effect of carbonaceous contaminants on the surface completely covering the support. Nevertheless, in both samples Au is still clearly visible at the outermost surface.

Based on the XPS and LEIS measurements, a segregation process upon heating in O<sub>2</sub>/CO is changing the surface composition of the Au/Ti nanoclusters, but in line with the computational investigation (see below) Au atoms are still present on the surface, thus opening up the possibility of Au catalysis using the Au/Ti clusters under realistic reaction conditions.

### 3.2 STEM study of the sintering process

The mean diameters of Au<sub>2057</sub> and Au/Ti (400k amu) clusters deposited onto thin silica film were measured, based on projected surface area, and found to be 4.07 nm and 3.94 nm, respectively. The shape of Au<sub>2057</sub> clusters can be obtained from the relationship between the cluster diameter ( $D$ ) and the number of Au atoms ( $N$ ).<sup>61</sup> The shape of Au cluster would be pseudospherical, if their  $D \sim N^{1/3}$  relations correspond with the spherical geometric model,  $D_s = \kappa N^{1/3}$ . Here,  $\kappa = 0.328$  nm, is calculated from the experimental data.<sup>62</sup> This value is corresponding well with twice of the Wigner–Seitz radius,  $r_{ws} = 0.165$  nm.<sup>63,64</sup> The diameters of the Au<sub>2057</sub> clusters fit a quasi-spherical geometric model well, indicating that they do not relax substantially to a hemispherical shape on the surface.

To evaluate whether the addition of Ti enhances the stability of the alloy clusters compared with the pure Au clusters, and to shed light on the sintering mechanism(s), the process of sintering induced by electron beam irradiation was investi-

gated with the aberration-corrected STEM. Fig. 3 shows sequential images of Au and Au/Ti cluster dimers exposed to the electron beam.

The observed sintering process of the Au clusters can be divided into two phases according to the STEM images and the change of the measured major axis of the Au dimer. In the first phase, the two Au clusters are seen to move toward each other once they are exposed to the electron beam driven by surface plasmon coupling,<sup>65,66</sup> which can be confirmed by the major axis of the Au dimer shrinking from 8.31 nm to 7.66 nm after 39 s, as plotted in Fig. 4 (the cluster migration can also be seen in the ESI in Fig. S2†). The cluster migration leads the two initially separated clusters to make contact with each other. Afterwards, in the second phase, a process of coalescence, presumably driven by peripheral atom diffusion,

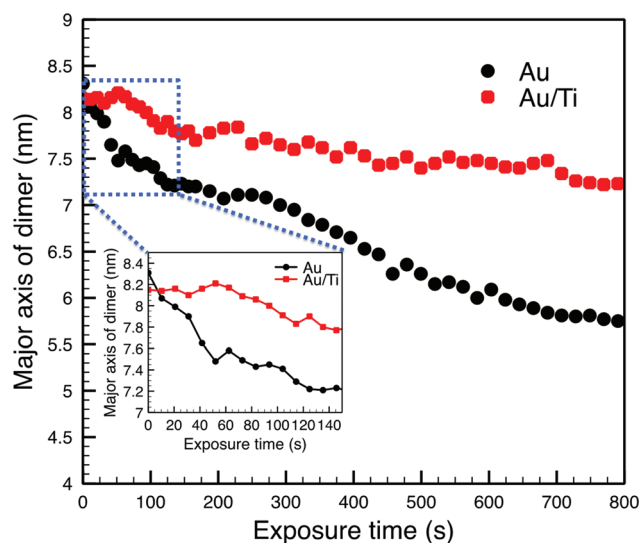


Fig. 4 The length of the major axis of the cluster dimers as a function of the time of exposure to the electron beam. The inset shows an enlarged view of the graph in blue square.

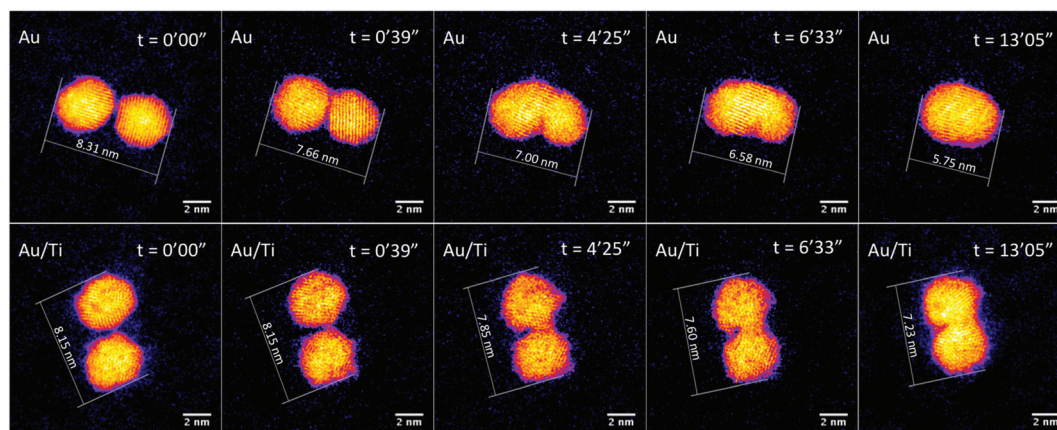


Fig. 3 STEM images of dimers of Au<sub>2057</sub> and of Au/Ti clusters continuously exposed to electron beam irradiation, with an acquisition time of 1.3 seconds per frame and a dose of  $6.3 \times 10^3$  e<sup>-</sup> per Ångström<sup>2</sup> per frame. At the beginning of imaging, the gap sizes between the Au<sub>2057</sub> and Au/Ti cluster pairs are ~0.30 nm and ~0.25 nm, respectively.



decreases the surface area of the dimer. Two different shrink rates of the major axis of the dimer can be found in these two phases of sintering, which are  $\sim 1.00 \text{ nm min}^{-1}$  and  $\sim 0.17 \text{ nm min}^{-1}$  for the first and second phases, respectively (Fig. 4). Thus the sintering behavior of the Au dimer starts at a relatively high rate with the two clusters quickly migrating towards each other. After about 50 s, once the clusters collide, the second and slower phase of sintering begins and the rate of peripheral Au atom diffusion regulates the rate of sintering. In the case of the Au<sub>2057</sub> clusters, 8 pairs of Au cluster dimers were imaged with the STEM, and 6 pairs of them showed similar sintering behavior as that discussed above. The clusters in the other 2 pairs of dimers with a gap size of  $\sim 0.8 \text{ nm}$  in between were simply found to move away from each other during imaging.

Compared with the Au clusters, the Au/Ti clusters show different sintering behavior, although once more it can be divided into two phases. In the first phase, the Au/Ti clusters remain at their original positions instead of quickly migrating towards each other like the Au clusters. Thus the Au/Ti cluster dimer retains the same length of major axis (8.15 nm) after being exposed to the electron beam for 39 s. In this first phase, only peripheral atom diffusion is found. However, once diffusing atoms found in the gap between the two clusters make a first physical connection between the clusters, the second phase of sintering is initiated. In this phase, coalescence takes place around the bridge formed in between, but the coalescence rate here ( $0.08 \text{ nm min}^{-1}$ ) is much lower than that of the Au cluster dimer, Fig. 4 (also be confirmed in Fig. S3 and S5 in the ESI†).

The different sintering behavior of Au/Ti clusters compared with Au clusters on silica can be attributed to the strong interaction between Ti and the lattice oxygen of the silica support, which tends to anchor the clusters in the dimer against sintering. However, the anchoring effect is insufficient to prevent

coalescence once two clusters come into contact, so it can slow down the coalescence process but cannot prevent it completely. If the distance between two neighbouring clusters is large enough (see below) that the diffusing atoms cannot “build a bridge” between them, then sintering is exceptionally slow.

Fig. 5 (see also Fig. S4 in the ESI†) shows a Au/Ti cluster dimer with a slightly larger gap than that in Fig. 3. It can be seen that the two clusters do not sinter after electron exposure for as long as 13 min, and basically the major axis of this dimer retains the same value. Peripheral atom diffusion is again found during electron exposure. At 1 min 28 s, the atoms highlighted by arrows can be found located between the two clusters, and then disappear 3 s later by moving away from the dimer area or binding to one of the two clusters. It is further observed that the cluster on the right exhibits a transient asperity due to the accumulation of diffusing atoms at 3 min 33 s. However, the distance between the clusters is too large for a contact bridge to be formed, and the asperity has decayed at 3 min 57 s. This confirms that sintering does not happen between Au/Ti clusters if the cluster distance is too large for the diffusing atoms to build a bridge in between. In future, experimental studies of hundreds of such dimer pairs may lead to the precise measurement of a critical distance for sintering by bridge formation, dependent on temperature and other key parameters (including in the case of the electron beam experiments, the current). In the case of Au/Ti clusters, 14 pairs of Au/Ti clusters dimers were imaged with STEM, and all of them showed similar behavior against sintering.

### 3.3 DFT calculations of surface anchoring

The size of the observed particles (around 2000 atoms) is such that we cannot directly model them with DFT calculations. It is likely that these particles will adopt shapes derived from a Wulff construction, such as a truncated octahedron dominated by (100) and (111) surfaces. To investigate further the role of Ti

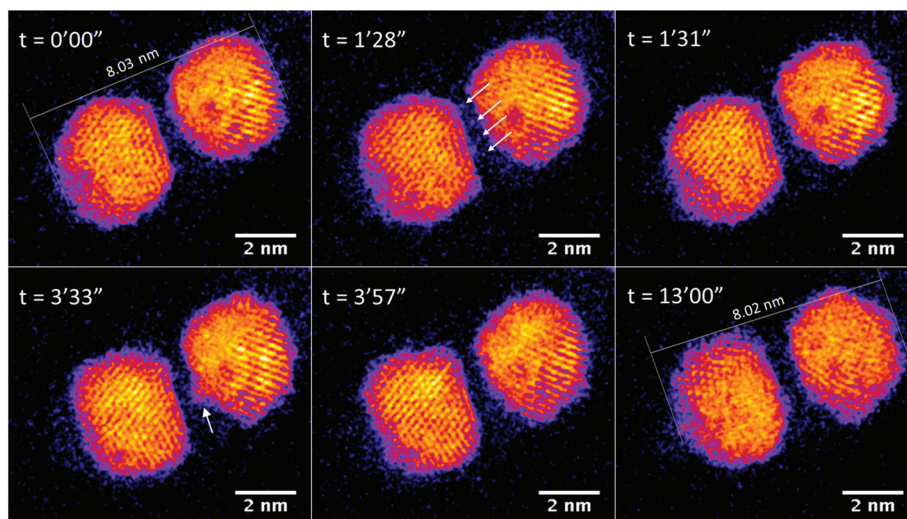


Fig. 5 STEM images of Au/Ti cluster dimers, with a larger gap ( $\sim 0.65 \text{ nm}$ ) between than Fig. 3, continuously exposed to electron beam irradiation with acquisition time of 1.3 seconds per frame and a dose of  $9.8 \times 10^3 \text{ e}^-$  per Angstrom<sup>2</sup> per frame.



atoms against sintering, we considered, amongst others, the adsorption of Au<sub>20</sub> and Au<sub>10</sub>Ti<sub>10</sub> tetrahedral clusters which exhibit (111) faces only. While these clusters are smaller than in the experiment, the nature of the bonding with the surface is rather local, and the essence of the interaction between cluster and support is sufficiently well represented by the chosen model. Before adsorbing the clusters on the surface, we optimized the geometries of the free-standing clusters. With photoelectron spectroscopy and relativistic density functional theory simulations, Li *et al.*<sup>67</sup> proposed that Au<sub>20</sub> prefers a tetrahedral shape. Wang and Palmer<sup>68</sup> found that this tetrahedral structure is observed, amongst other isomers, when size-selected Au<sub>20</sub> clusters soft-landed were observed by aberration-corrected STEM on amorphous carbon film. Other experimental<sup>69</sup> and computational<sup>70</sup> studies also identify the tetrahedral geometry. In our simulations we also find that the tetrahedral structure is the lowest energy geometry for free Au<sub>20</sub> clusters. For the bimetallic Au<sub>10</sub>Ti<sub>10</sub> cluster, we considered this and other structures, while also testing different distributions of the two elements within the clusters. We considered several cluster isomers, but the potential energy surface (PES) possesses a very large number of local minima, which cannot be extensively explored without a global search algorithm. In this respect, our results are representative of some potential structures, but we do not pretend to have identified the global minimum on the PES. We find that the best cluster shape for the Au<sub>10</sub>Ti<sub>10</sub> cluster is not in fact a pyramid, but a compact shape with Ti atoms sitting largely inside the cluster and a majority of Au atoms in the outer layer.

After the optimization of the free-standing clusters, the most stable cluster isomers were deposited on the fully hydroxylated  $\alpha$ -quartz (001) surface. This choice is dictated by the assumption that under experimental conditions the silica surface is not hydroxyl-free. Fully dehydroxylated silica surfaces can be obtained only after thermal treatment above 600 °C,<sup>71</sup> and SiO<sub>2</sub> surfaces get partially hydroxylated even under high vacuum conditions almost instantly.<sup>72</sup> Since our samples are exposed to air, it is likely that a given concentration of OH groups will be present. The density of the OH groups in our model is probably higher than in the real samples, but no quantitative assessment is possible. In any case, in a previous study on tiny AuTi bimetallic clusters we considered both hydroxylated and de-hydroxylated silica surfaces, and we found that the AuTi bimetallic clusters interact strongly also with the de-hydroxylated surface. In this case oxygen atoms are extracted from the surface to bind strongly with the Ti component of the cluster.<sup>73</sup> The structure of the hydroxylated silica surface has been described in a previous study.<sup>27</sup> The resulting structures are shown in Fig. 6, and relevant adsorption parameters are summarized in Table 1.

We find that the Au<sub>20</sub> cluster exhibits a relatively low adsorption energy on the silica surface, only -1.05 eV, Table 1. The atoms in the supported Au clusters show an average Bader charge close to zero, indicating the absence of chemical interaction with the surface. The relatively weak binding of the Au cluster on the surface is thus largely due to dispersion forces.

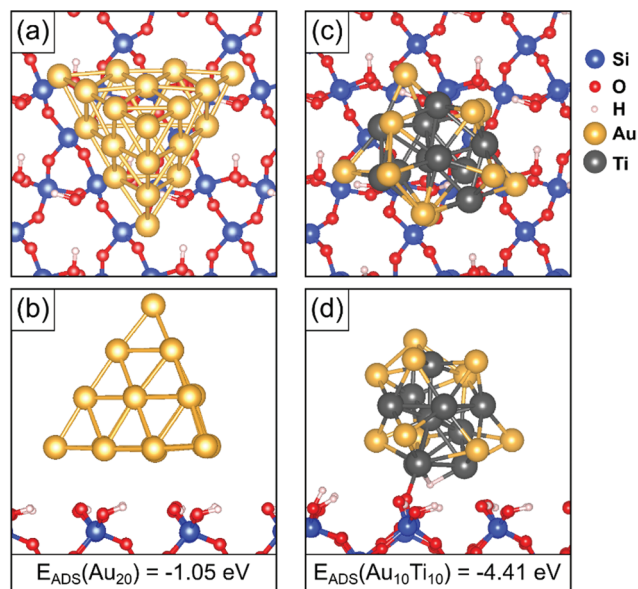


Fig. 6 (a) Top and (b) side views of Au<sub>20</sub> supported on the fully hydroxylated  $\alpha$ -quartz (001) surface from DFT simulations. (c) Top view and (d) side view of Au<sub>10</sub>Ti<sub>10</sub> supported on the same surface.

Table 1 Adsorption energies  $E_{\text{ADS}}$  (eV), average Bader charges  $q_{\text{AVG}}$  (Au or Ti) ( $|e|$ ) on the Au or Ti atoms in the cluster and total Bader charge on the cluster  $q_{\text{TOT}}(\text{cluster})$  ( $|e|$ ) from DFT simulations

	$E_{\text{ADS}}$	$q_{\text{AVG}}(\text{Au})$	$q_{\text{AVG}}(\text{Ti})$	$q_{\text{TOT}}(\text{cluster})$
Au <sub>20</sub> free-standing	—	0.00	—	0.00
Au <sub>20</sub> on silica	-1.05	0.00	—	0.00
Au <sub>10</sub> Ti <sub>10</sub> free-standing	—	-0.77	+0.77	0.00
Au <sub>10</sub> Ti <sub>10</sub> on silica	-4.41	-0.79	+0.93	+1.39

The Au<sub>10</sub>Ti<sub>10</sub> bimetallic cluster, on the other hand, exhibits a much larger adsorption energy of -4.41 eV. Only around 20% of the adsorption energy is due to dispersion forces. The large adsorption energy is due to a spontaneous local reaction of the cluster with the surface, *via* hydrogen reverse spillover. This can be seen in the side view in Fig. 6(d). The surface  $\equiv\text{Si}-\text{O}-\text{H}$  group is split, enabling the H atom to bind to the cluster at a Ti hollow site, while the residual  $\equiv\text{Si}-\text{O}^{\bullet}$  group binds to an adjacent Ti atom of the cluster, with a resulting Ti-O bond length of 1.96 Å. This local reaction anchors the cluster to the support *via* the formation of a  $\equiv\text{Si}-\text{O}-\text{Au}_{10}\text{Ti}_{10}$  complex. We expect this strong local binding interaction to anchor the AuTi bimetallic clusters to the hydroxylated silica surface, and to operate independently of the cluster size, leading to reduced sintering, as observed in the experiments.

Of course, in the experimental situation the clusters are exposed to ambient conditions between deposition and the sintering experiments. Thus we may expect some oxidation of the Ti atoms at the periphery of the Au/Ti clusters, after the removal of the sample out of the deposition chamber, as the XPS confirms that Ti is present on the surface in the form of



TiO<sub>2</sub>. However, during and directly after the deposition of the clusters in UHV, the metallic clusters can react with the silica surface, as shown in the calculations.

## 4. Conclusions

Size-selected Au/Ti nanoalloy clusters (400, 000 amu) and pure Au<sub>2057</sub> (405, 229 amu) clusters were produced by gas-phase synthesis with a magnetron sputtering, gas condensation cluster beam source and deposited onto silica supports. Chemical characterisation of the deposited clusters was provided by XPS and LEIS measurements. Upon heating the clusters in an O<sub>2</sub>/CO environment to simulate catalytic reaction conditions, the surface composition of the Au/Ti clusters was changed by a segregation process, but both Au and Ti were still visible on the surface. Thus the supported nanoalloy Au/Ti clusters are considered as viable candidates for Au-based catalysis.

The sintering behavior of the clusters under electron beam annealing was explored by aberration-corrected STEM imaging in real space and real time. Two neighbouring Au<sub>2057</sub> clusters in a dimer were found to quickly migrate and, in a second slower phase, coalesce with each other. In contrast, Au/Ti dimers showed a strong anchoring effect against sintering due to the presence of the reactive Ti atoms, most notably when the gap between them exceeded 0.60 nm. Sintering can still happen if two Au/Ti clusters are extremely close to each other, but this is due to atom diffusion between the two clusters instead of cluster migration. Sintering is expected to be exceedingly slow if the distance between the Au/Ti clusters is large enough and the diffusing atoms cannot “build a bridge” in between.

DFT calculations show that, in model bimetallic clusters (20 atoms), the Au atoms prefer a position at the surface of the nanoclusters, in good agreement with the outcome of the LEIS experiments. The calculations furthermore show that Au clusters can be bonded much more strongly (by a factor of 5) to a silica support by alloying them with Ti. This effect is due to the increased reactivity of the Au/Ti bimetallic clusters when they present a surface containing a reactive metal (Ti). Future computations may address the role of the oxidation of the bimetallic Au/Ti clusters but the qualitative enhancement of surface anchoring against sintering is expected to be preserved. The presence of surface Au atoms in size-selected Au/Ti nanoalloy clusters and the anchoring effect due to Ti incorporation may open up new possibilities in Au-based nanocatalysis.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

For funding of this work, we gratefully acknowledge the European Union's Seventh Framework Programme (FP7/2007-2013) under grant agreement no. 607417 (Catsense), as well as

an EPSRC Fellowship Grant and the Villum Foundation V-SUSTAIN grant 9455 to the Villum Center for the Science of Sustainable Fuels and Chemicals.

## References

- 1 B. Hammer and J. K. Nørskov, *Nature*, 1995, **376**, 238–240.
- 2 W. A. Bone and R. V. Wheeler, *Proc. R. Soc. A*, 1906, **77**, 146–147.
- 3 W. A. Bone and G. W. Andrew, *Proc. R. Soc. A*, 1925, **109**, 459–476.
- 4 K. Sennewald, W. Vogt and H. Glaser, *Patent* DE1244766B, 1967.
- 5 G. C. Bond and P. A. Sermon, *Gold Bull.*, 1973, **6**, 102–105.
- 6 M. Haruta, T. Kobayashi, H. Sano and N. Yamada, *Chem. Lett.*, 1987, 405–408.
- 7 M. Haruta, N. Yamada, T. Kobayashi and S. Iijima, *J. Catal.*, 1989, **115**, 301–309.
- 8 G. Hutchings, *J. Catal.*, 1985, **96**, 292–295.
- 9 M. Haruta, *Angew. Chem., Int. Ed.*, 2014, **53**, 52–56.
- 10 A. Taketoshi and M. Haruta, *Chem. Lett.*, 2014, **43**, 380–387.
- 11 B. Hvolbæk, T. V. W. Janssens, B. S. Clausen, H. Falsig, C. H. Christensen and J. K. Nørskov, *Nano Today*, 2007, **2**, 14–18.
- 12 Y. G. Wang, D. C. Cantu, M. S. Lee, J. Li, V. A. Glezakou and R. Rousseau, *J. Am. Chem. Soc.*, 2016, **138**, 10467–10476.
- 13 J. Saavedra, C. Powell, B. Panthi, C. J. Pursell and B. D. Chandler, *J. Catal.*, 2013, **307**, 37–47.
- 14 K. Christmann, S. Schwede, S. Schubert and W. Kudernatsch, *ChemPhysChem*, 2010, **11**, 1344–1363.
- 15 B. K. Min and C. M. Friend, *Chem. Rev.*, 2007, **107**, 2709–2724.
- 16 D. Widmann and R. J. Behm, *Acc. Chem. Res.*, 2014, **47**, 740–749.
- 17 T. Fujitani and I. Nakamura, *Angew. Chem., Int. Ed.*, 2011, **50**, 10144–10147.
- 18 I. N. Remediakis, N. Lopez and J. K. Nørskov, *Angew. Chem., Int. Ed.*, 2005, **117**, 1858–1860.
- 19 Z. Duan and G. Henkelman, *ACS Catal.*, 2015, **5**, 1589–1595.
- 20 I. X. Green, W. J. Tang, M. Neurock and J. T. Yates, *Acc. Chem. Res.*, 2014, **47**, 805–815.
- 21 I. X. Green, W. J. Tang, M. Neurock and J. T. Yates, *Science*, 2011, **333**, 736–739.
- 22 P. Kast, G. Kučerová and R. J. Behm, *Catal. Today*, 2015, **244**, 146–160.
- 23 D. Widmann and R. J. Behm, *Angew. Chem., Int. Ed.*, 2011, **50**, 10241–10245.
- 24 L. Zhang, R. Persaud and T. E. Madey, *Phys. Rev. B: Condens. Matter*, 1997, **56**, 10549–10557.
- 25 S. C. Parker, A. W. Grant, V. A. Bondzie and C. T. Campbell, *Surf. Sci.*, 1999, **441**, 10–20.
- 26 G. M. Veith, A. R. Lupini, S. Rashkeev, S. J. Pennycook, D. R. Mullins, V. Schwartz, C. A. Bridges and N. J. Dudney, *J. Catal.*, 2009, **262**, 92–101.



- 27 P. Schlexer and G. Pacchioni, *Top. Catal.*, 2016, **60**, 459–470.
- 28 F. Yang, M. S. Chen and D. W. Goodman, *J. Phys. Chem. C*, 2009, **113**, 254–260.
- 29 K. J. Hu, S. R. Plant, P. R. Ellis, C. M. Brown, P. T. Bishop and R. E. Palmer, *J. Am. Chem. Soc.*, 2015, **137**, 15161–15168.
- 30 A. Cao, R. Lu and G. Veser, *Phys. Chem. Chem. Phys.*, 2010, **12**, 13499–13510.
- 31 Z. Ma and S. Dai, *Heterogeneous Gold Catalysts and Catalysis*, Royal Soc Chemistry, Cambridge, 2014.
- 32 P. Hernandez-Fernandez, F. Masini, D. N. McCarthy, C. E. Strebel, D. Friebel, D. Deiana, P. Malacrida, A. Nierhoff, A. Bodin, A. M. Wise, J. H. Nielsen, T. W. Hansen, A. Nilsson, I. E. L. Stephens and I. Chorkendorff, *Nat. Chem.*, 2014, **6**, 732–738.
- 33 A. Velazquez-Palenzuela, F. Masini, A. F. Pedersen, M. Escudero-Escribano, D. Deiana, P. Malacrida, T. W. Hansen, D. Friebel, A. Nilsson, I. E. L. Stephens and I. Chorkendorff, *J. Catal.*, 2015, **328**, 297–307.
- 34 C. Holve, C. F. Elkjaer, A. Nierhoff, J. Sehested, I. Chorkendorff, S. Helveg and J. H. Nielsen, *J. Phys. Chem. C*, 2015, **119**, 2804–2812.
- 35 A. Gomez-Cortes, G. Diaz, R. Zanella, H. Ramirez, P. Santiago and J. M. Saniger, *J. Phys. Chem. C*, 2009, **113**, 9710–9720.
- 36 Y. J. Guan and E. J. M. Hensen, *J. Catal.*, 2013, **305**, 135–145.
- 37 C. W. Han, P. Majumdar, E. E. Marinero, A. Aguilar-Tapia, R. Zanella, J. Greeley and V. Ortalan, *Nano Lett.*, 2015, **15**, 8141–8147.
- 38 X. Y. Liu, A. Q. Wang, X. D. Wang, C. Y. Mou and T. Zhang, *Chem. Commun.*, 2008, 3187–3189.
- 39 X. Y. Liu, A. Q. Wang, X. F. Yang, T. Zhang, C. Y. Mou, D. S. Su and J. Li, *Chem. Mater.*, 2009, **21**, 410–418.
- 40 G. C. Ma, A. Binder, M. F. Chi, C. Liu, R. C. Jin, D. E. Jiang, J. Fan and S. Dai, *Chem. Commun.*, 2012, **48**, 11413–11415.
- 41 A. Sandoval, C. Louis and R. Zanella, *Appl. Catal., B*, 2013, **140**, 363–377.
- 42 B. von Issendorff and R. E. Palmer, *Rev. Sci. Instrum.*, 1999, **70**, 4497.
- 43 S. Pratontep, S. J. Carroll, C. Xirouchaki, M. Streun and R. E. Palmer, *Rev. Sci. Instrum.*, 2005, **76**, 045103.
- 44 G. Kresse and J. Furthmuller, *Phys. Rev. B: Condens. Matter*, 1996, **54**, 11169–11186.
- 45 G. Kresse and J. Furthmuller, *Comput. Mater. Sci.*, 1996, **6**, 15–50.
- 46 G. Kresse and J. Hafner, *Phys. Rev. B: Condens. Matter*, 1993, **47**, 558–561.
- 47 G. Kresse and J. Hafner, *Phys. Rev. B: Condens. Matter*, 1994, **49**, 14251–14269.
- 48 J. P. Perdew, K. Burke and M. Ernzerhof, *Phys. Rev. Lett.*, 1996, **77**, 3865–3868.
- 49 J. P. Perdew, K. Burke and M. Ernzerhof, *Phys. Rev. Lett.*, 1997, **78**, 1396–1396.
- 50 P. E. Blochl, *Phys. Rev. B: Condens. Matter*, 1994, **50**, 17953–17979.
- 51 G. Kresse and D. Joubert, *Phys. Rev. B: Condens. Matter*, 1999, **59**, 1758–1775.
- 52 G. H. F. Diercksen and S. Wilson, *Methods in Computational Molecular Physics*, Springer Netherlands, 1983.
- 53 C. Moler and I. Shavitt, *Report on the Workshop: Numerical Algorithms in Chemistry: Algebraic Methods*, Lawrence Berkeley Laboratory, University of California, 1978.
- 54 P. Schlexer and G. Pacchioni, *J. Phys.: Condens. Matter*, 2016, **28**, 364005.
- 55 H. J. Monkhorst and J. D. Pack, *Phys. Rev. B: Condens. Matter*, 1976, **13**, 5188–5192.
- 56 S. Grimme, *J. Comput. Chem.*, 2006, **27**, 1787–1799.
- 57 S. Tosoni and J. Sauer, *Phys. Chem. Chem. Phys.*, 2010, **12**, 14330–14340.
- 58 G. Henkelman, A. Arnaldsson and H. Jonsson, *Comput. Mater. Sci.*, 2006, **36**, 354–360.
- 59 E. Sanville, S. D. Kenny, R. Smith and G. Henkelman, *J. Comput. Chem.*, 2007, **28**, 899–908.
- 60 W. Tang, E. Sanville and G. Henkelman, *J. Phys.: Condens. Matter*, 2009, **21**, 084204.
- 61 N. Young, Z. Li, Y. Chen, S. Palomba, M. Di Vece and R. Palmer, *Phys. Rev. Lett.*, 2008, **101**, 246103.
- 62 K. P. Arkill, J. M. Mantell, S. R. Plant, P. Verkade and R. E. Palmer, *Sci. Rep.*, 2015, **5**, 9234.
- 63 Y. Han, D. S. He, Y. Liu, S. Xie, T. Tsukuda and Z. Y. Li, *Small*, 2012, **8**, 2361–2364.
- 64 B. Smirnov, *Clusters and small particles: in gases and plasmas*, Springer, New York, 2000.
- 65 P. E. Batson, *Phys. Rev. Lett.*, 1982, **49**, 936–940.
- 66 P. E. Batson, A. Reyes-Coronado, R. G. Barrera, A. Rivacoba, P. M. Echenique and J. Aizpurua, *Nano Lett.*, 2011, **11**, 3388–3393.
- 67 J. Li, X. Li, H. J. Zhai and L. S. Wang, *Science*, 2003, **299**, 864–867.
- 68 Z. W. Wang and R. E. Palmer, *Nanoscale*, 2012, **4**, 4947–4949.
- 69 P. Gruene, D. M. Rayner, B. Redlich, A. F. G. van der Meer, J. T. Lyon, G. Meijer and A. Fielicke, *Science*, 2008, **321**, 674–676.
- 70 L. M. Molina and B. Hammer, *J. Catal.*, 2005, **233**, 399–404.
- 71 C. Armistead, A. Tyler, F. Hambleton, S. Mitchell and J. A. Hockey, *J. Phys. Chem.*, 1969, **73**, 3947–3953.
- 72 A. S. D'Souza and C. G. Pantano, *J. Am. Ceram. Soc.*, 1999, **82**, 1289–1293.
- 73 P. Schlexer and G. Pacchioni, *J. Phys. Chem. C*, 2017, 14717–14724.

